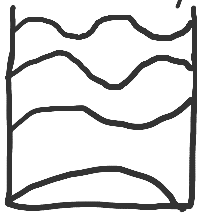


Fermi and Bose distributions

Fermions - in each quantum state there cannot simultaneously be > 1 particle
= Fermi-Dirac statistics

Bosons - in each quantum state there may be > 1 particles

Consider a system of non-interacting particles in with quantum states labelled by k



← levels of one particle

In principle, each of those levels may be considered as a separate subsystem

Indeed,

$$\Omega = -T \ln \sum_N e^{\frac{\mu N}{T}} \sum_i e^{-\frac{E_{iN}}{T}}$$

$$N = n_1 + n_2 + \dots$$

$$E_{nN} = n_1 \epsilon_1 + n_2 \epsilon_2 + \dots$$

$$\Omega = -T \ln \sum_{\{n\}} e^{\frac{\mu}{T}(n_1 + n_2 + \dots)} e^{-\frac{\epsilon_1 n_1 + \epsilon_2 n_2 + \dots}{T}} =$$

$$= -T \ln \sum_{\{n\}} \prod_k e^{\frac{\mu - \epsilon_k}{T} n_k} =$$

$$= -T \ln \sum_{\{n_k\}} e^{-\frac{\mu - \epsilon_k}{T} n_k} = \sum_k \Omega_k$$

where $\Omega_k = -T \ln \sum_{n_k} e^{-\frac{\mu - \epsilon_k}{T} n_k}$

For fermions $n_k = 0$ or $n_k = 1$

$$\text{Then } \Omega_k = -T \ln (1 + e^{-\frac{\mu - \epsilon_k}{T}})$$

The average number of fermions in a state labelled by k is

$$\bar{n}_k \equiv -\frac{\partial \Omega_k}{\partial \mu} = \frac{e^{-\frac{\mu - \epsilon_k}{T}}}{1 + e^{-\frac{\mu - \epsilon_k}{T}}} = \frac{1}{e^{\frac{\epsilon_k - \mu}{T}} + 1}$$

$$f = \frac{1}{e^{\frac{\epsilon - \mu}{T}} + 1}$$

- Fermi-Dirac aka Fermi distribution

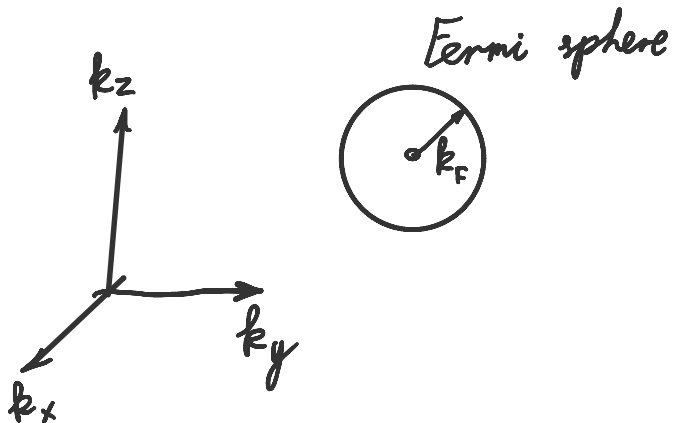
$$0 \leq f \leq 1$$

One should use the normalisation condition $\sum_k n_k = N$
 For large temperatures $T \gg \mu$ $n_k \approx e^{-\frac{\epsilon_k - \mu}{T}}$

For a given chemical potential μ
 low energies $\epsilon \rightarrow -\infty$ correspond to $f=1$ (filled states)
 High-energy states with $\epsilon \rightarrow +\infty$ are occupied
 + D.

High-energy states with $\epsilon \gg k_B T$

Example: electrons in a metal:



Bosons

$$\Omega_k = -T \ln \sum_{n_k=0}^{\infty} e^{\frac{\mu - \epsilon_k}{T} n_k}$$

The geometric series is convergent only for $e^{\frac{\mu - \epsilon_k}{T}} < 1 \iff \epsilon_k > \mu$

If $\epsilon_k = 0$ is the ground state, as it is often convenient to choose,

$$\Omega_k = -T \ln \frac{1}{1 - e^{\frac{\mu - \epsilon_k}{T}}} = T \ln \left(1 - e^{\frac{\mu - \epsilon_k}{T}} \right)$$

$$\bar{n}_k \equiv - \frac{\partial \Omega_k}{\partial \mu} = \frac{1}{e^{\frac{\epsilon_k - \mu}{T}} - 1}$$

which

$$n_k = \frac{1}{e^{\frac{\epsilon_k - \mu}{T}} - 1}$$

- Bose-Einstein
aka Bose
distrib -

Example: phonons in an oscillator: $\frac{1}{e^{\frac{\hbar\omega}{T}} - 1}$